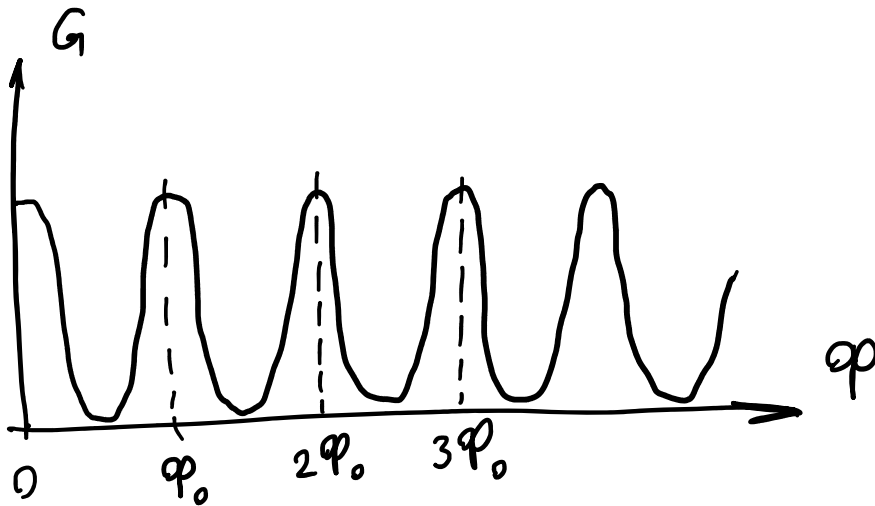


Now the wave function

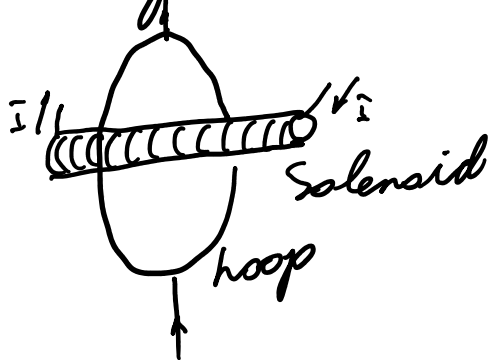
$$\Psi = \Psi_L + \Psi_R \sim e^{-i\pi\frac{\varphi}{\varphi_0}} + e^{i\pi\frac{\varphi}{\varphi_0}} = \cos\left(\pi\frac{\varphi}{\varphi_0}\right)$$

$$|\Psi|^2 \sim \cos^2\left(\pi\frac{\varphi}{\varphi_0}\right) \sim 1 + \cos\left(2\pi\frac{\varphi}{\varphi_0}\right)$$

Conductance



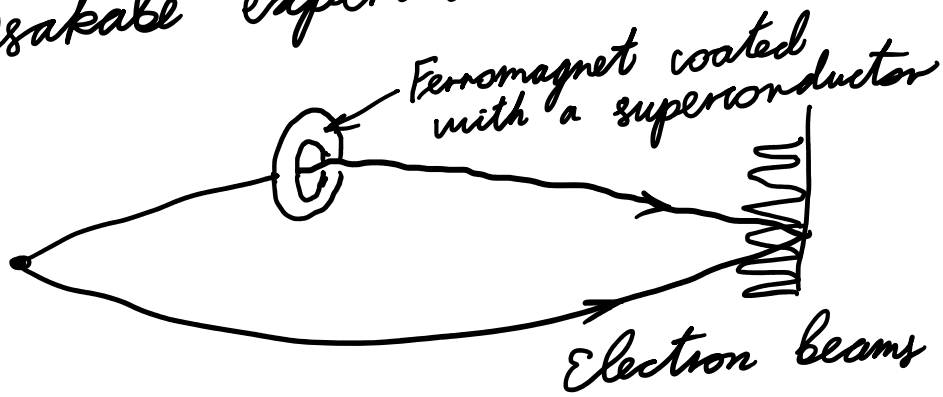
Note: Aharonov-Bohm oscillations exist even when the magnetic field does not reach the wires; it may be confined to the solenoid which is threading the wire loop



$$\oint \vec{A} \cdot d\vec{r} = \varphi$$

Nakabe experiment: (1986)

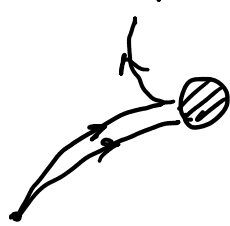
Osakabe experiment: (1980)



Effect of magnetic field on WL

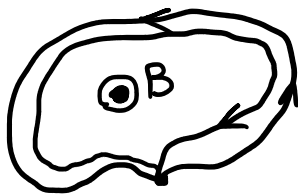
Effect of magnetic field on transport:
magnetic field makes trajectories curve

$$\omega_c = \frac{eB}{mc}$$



- cyclotron frequency
If $\omega_c \tau \ll 1$, then there is another much stronger effect: effect on WL

WL comes from the interference of pairs of trajectories propagating in opposite directions



$$|A_1 + A_2|^2 =$$

$$= |A_1|^2 + |A_2|^2 + \underbrace{A_1^* A_2 + A_2^* A_1}_{2 \cos(\pi \frac{\phi}{\phi_0})}$$

$$\langle \cos(\pi \frac{\phi}{\phi_0}) \rangle = 0$$

$$\left\langle \cos\left(\pi \frac{\varphi}{\varphi_0}\right) \right\rangle_{\text{all trajectories}} = 0$$

→ Magnetic field destroys WL

Focus on 2D

The characteristic "interference area" is

$$S_{\varphi} = D \tau_{\varphi}$$

$$\rightarrow B_{\varphi} = \frac{c \hbar}{e} \frac{1}{D \tau_{\varphi}}$$

If we write $D \sim \frac{E_F \tau}{m}$ then

$$\underbrace{\frac{e B_{\varphi} \tau}{m c}}_{\omega_c} \underbrace{\frac{\tau_{\varphi} E_F}{\hbar}}_{\approx 1} \approx 1 \rightarrow \omega_c \tau \underbrace{\frac{E_F \tau_{\varphi}}{\hbar}}_{\approx 1 \text{ in a good metal}} \approx 1$$

$$\rightarrow \omega_c \tau \ll 1$$

